Solar Cycle Prediction and Reconstruction

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Outline

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Polar magnetic fields – producing the next cycle
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Solar Cycle Characteristics
Cycle Variability

Yearly average of daily sunspot numbers (V2.0) since the invention of the telescope in 1610.

- Cycle period ~11 years (9-14 years)
- Cycle amplitude ~ 150 (50-250)
- Maunder Minimum (no sunspots 1645-1715)
- Gleissberg Cycle (~100 year modulation of cycle amplitudes)
The average sunspot cycle (red curve) is asymmetric – rising to a maximum in 4 years and declining to minimum over the remaining 7 years. This shape is well fit with a parametric function:

\[ F(t) = A \left( \frac{t - t_0}{b} \right)^3 \left[ \exp \left( \frac{t - t_0}{b} \right)^2 - c \right]^{-1} \]

with \( c=0.71 \) and \( b=b(A) \) leaving a function of just \( t_0 \) and \( A \) (Hathaway et al. 1994).
Predicting Activity in an Ongoing Cycle

Reliable predictions of ongoing cycles have been made by fitting the observed monthly averaged sunspot numbers to this parametric function. These predictions become reliable 2-3 years after the cycle starts (blue dots at inflection points of curve).
Producing the Cycle – the Solar Dynamo
The Babcock-Leighton Dynamo

a. Cycle starts with dipole field at minimum

b. Dipole field is stretched by differential rotation to produce azimuthal, toroidal field

c. Toroidal field becomes strong enough to become buoyant and rise through the surface to form bipolar active regions with a tilt from the azimuthal direction

d. Surface flux transport carries high-latitude (predominantly) following polarity flux to the poles

a. New (reversed polarity) polar fields are the seeds of the next cycle

Babcock (1961)
The Dynamo Observed

The radial component of the photospheric magnetic field averaged over longitude for each 27-day rotation of the Sun over the last four cycles.

- Active latitude butterfly wings have following polarity at high latitudes and preceding polarity at low latitudes.
- Poleward transport of the predominantly following polarity flux reverses the polar fields at maximum and builds up new polar fields by the next minimum.
Polar Fields – Producing the next Cycle
The average polar field strength at cycle minimum (blue bars) correlates well with the maximum sunspot number of next cycle (black arrows).
Geomagnetic activity at cycle minima (black dots) is produced by high-speed solar wind streams from coronal holes that carry magnetic fields representative of the dipole field.

The minima in the geomagnetic aa-index are well correlated with the maximum sunspot number of the next cycle (Ohl, 1966).

Polar faculae (produced by polar magnetic elements) are well correlated with the maximum sunspot number of the next cycle (Muñoz-Jaramillo et al., 2012).
Surface Flux Transport – Producing the Polar Fields
Active regions emerge with a characteristic tilt that leaves the following-polarity flux at higher latitudes.

As the cycle progresses active regions emerge at lower and lower latitudes where preceding-polarity flux is canceled and replaced with following-polarity flux.

This high-latitude, following-polarity flux is transported to the poles by random convective motions and the poleward directed meridional flow.
Surface Flux Transport – Four Decades
Surface Flux Transport – Four Days
The Surface Flux Transport Promise

The evolution of the Sun’s surface magnetic field can be obtained -

1. Given the emergence of (tilted) active region magnetic flux
2. Given knowledge of the flows (the differential rotation, the meridional circulation, the cellular convective flows) and their variations

Can we use surface flux transport to determine what the strength of the polar fields will be at the cycle minimum well before the end of the cycle?
Measuring the Flux Transport Flows
Full disk magnetograms from SOHO/MDI and SDO/HMI are mapped into heliographic longitude and latitude. Long, narrow strips at each latitude are cross-correlated with similar strips obtained 8 hours later to find the shift in longitude and latitude that gives the best correlation (Hathaway & Rightmire 2010, 2011).
The Differential Rotation

The differential rotation (axisymmetric longitudinal flow) is faster than average at the equator and slower than average at higher latitudes.

We obtained hourly measurements of the latitudinal profiles, averaged them over 27-day solar rotations, and fit the profiles to 4th order polynomials.

Those profiles are well fit with three Legendre polynomials whose coefficients ($T_0$, $T_2$, and $T_4$) vary only slightly with cycle phase – faster and flatter at cycle maximum.

While differential rotation has a huge effect on the longitudinal structure, it does not impact the polar fields.
The Meridional Circulation

The meridional flow is very weak and has thus been difficult to measure.

The flow is poleward in each hemisphere and peaks at 30-40°.

Measurements with SDO/HMI up to 85° show flow to the poles without any evidence for counter-cells at high latitudes.

The meridional flow profiles are well fit by two Legendre polynomials with coefficients $S_1$ and $S_3$.

Our measurements (and earlier measurements by Komm, Howard, & Harvey, 1993) show that the flow is fast at cycle minima, slow at cycle maxima, and varies from cycle to cycle.

![Graph showing northward velocity and fit coefficients over time](image)
Deviations from the Average

Removing the average differential rotation profile from the profiles for each individual solar rotation reveals the “Torsional Oscillations” – faster rotation on the equatorward sides of the active latitude bands (red dots indicate sunspot area centroids) and slower rotation on the poleward sides.

The high-latitude spin-up around the time of cycle maximum produces the faster rotation and flatter rotation profiles at maxima.

Removing the average meridional flow profile from the profiles for each individual solar rotation reveals “Active Latitude Inflows” – poleward flow on the equatorward sides of the active latitudes and equatorward flows on the poleward sides.

These inflows produce the weaker meridional flow at cycle maxima and they vary with the level of activity in a cycle.
Hathaway et al. (2015) analyzed and simulated Doppler velocity data from SDO/HMI to determine the spectrum of flow velocities for all three vector velocity components – radial velocity, poloidal velocity (horizontal flows with divergence), and toroidal flows (horizontal flows with curl).

The total velocity increases with wavenumber to a peak of ~500 m/s for cells the size of supergranules (30 Mm in diameter) and then rises again to a peak of ~3000 m/s for cells the size of granules (1200 km in diameter).

Toroidal (heliostrophic?) flows dominate at wavenumbers less than 30.

There is only weak evidence for changes in the spectrum with the solar cycle.
“Giant Cells”

Hathaway, Upton, & Colegrove (2013) found that the largest cellular flow structures on the Sun had some expected flow properties – cyclonic flows with hemisphere dependent kinetic helicity and an equatorward flux of zonal momentum.
The Full Convection Spectrum

The velocity spectrum from our analysis of SDO/HMI Doppler velocity data agrees well with the spectrum of the “Giant Cells” from supergranule correlation tracking at the low wavenumber end (black line), and also agrees well with the spectrum of the flows in the StellarBox Radiative-MHD simulation of Wray et al. (2015) at the high wavenumber end (blue).
Reconstructing and Predicting the Cycle
We have developed an advective surface flux transport code that employs our measured surface flows to transport magnetic flux assimilated from either full-disk magnetic images of the Sun or from active region databases (as here). It reproduces the observed magnetic features – including the polar fields and their reversals. These simulations can be run into the future (using a statistical model of active region emergence) to predict the polar fields.
Testing the Flux Transport Promise

After 6 months

Full Disk Data Assimilation

After 1 year

Active Region Sources

After 2 years
Longitudinal Averages

Full-disk Data assimilation

Active Region Sources

Polar Field Strength (G)

2010  2015
The axial dipole moment measured by the Wilcox Solar Observatory shows the reversal in early 2013 and a rise to absolute values near those seen in early 2011 by early 2015.

Our (Upton & Hathaway, 2013) polar field prediction based on surface flux transport using active region sources from cycle 14 (a similar cycle) shows good agreement.

Note: Using different realizations of the convective motions (multi-colored tracks) gives diverging results - e.g. axial dipole = 0.8 ± 0.1 after 3.5 years.
Cycle 23 and Meridional Flow Variations

*Upton & Hathaway (2014)* showed that the observed variations in the meridional flow during cycle 23 produced **stronger** polar fields than would be produced with a constant, average, meridional flow profile.

This suggests that the changes to the meridional flow induced by the presence of active regions modulates the cycle amplitudes (*Cameron & Schüssler, 2012*).

Note: Changes in the meridional flow can produce significant changes in the results – e.g. axial dipole maximum = 1.4 ± 0.2.
Cameron et al. (2013) noted that the emergence of just a few badly oriented active regions can have significant consequences for the polar fields.

Note: Random perturbations in the tilt of large active regions can alter the axial dipole.
Conclusions

• We can predict the level of activity in an ongoing cycle once a cycle is well underway (curve-fitting)
• We can predict the amplitude and timing of a future cycle based on observed polar fields at the time of minimum
• We can predict polar fields at least 2-3 years ahead based on surface flux transport with active region sources based on statistics
• Long range predictions (more than one cycle ahead) are limited by:
  1. Random variations in the convective flows
  2. Future variations in the meridional flow
  3. Random variations in the tilt of active regions

Caveat Emptor